

UNIVERSITY OF SWAZILAND

FACULTY OF SCIENCE

DEPARTMENT OF PHYSICS

MAIN EXAMINATION 2009-10

TITLE OF THE PAPER: QUANTUM MECHANICS-I

COURSE NUMBER : P342

TIME ALLOWED : THREE HOURS

INSTRUCTIONS:

- ANSWER ANY **FOUR** OUT OF **FIVE** QUESTIONS.
- EACH QUESTION CARRIES **25** MARKS.
- MARKS FOR DIFFERENT SECTIONS ARE SHOWN IN THE RIGHT-HAND MARGIN.
- USE THE INFORMATION GIVEN IN THE ATTACHED **APPENDIX** WHEN NECESSARY.

THIS PAPER HAS **Six** PAGES, INCLUDING THIS PAGE.

DO NOT OPEN THE PAPER UNTIL THE INVIGILATOR HAS GIVEN PERMISSION.

Q.1. (a) What is the de-Broglie wavelength of neutron with a kinetic energy of 1.7 MeV? Will such neutrons produce a diffraction pattern in a crystal of lattice point distances of the order of 10^{-8} m? [3]

(b) Using the uncertainty relation, show that in a nucleus with average potential energy $\langle U \rangle \geq 15$ MeV, the bound nucleon is confined within a sphere of radius $r_0 \geq 1.2 \times 10^{-15}$ m. [2]

(c) Explain [6]

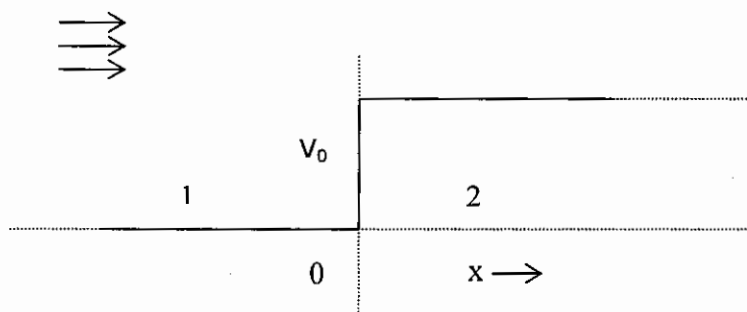
(i) Superposition principle and its importance in quantum mechanics.

(ii) Why it is necessary to have only linear operators associated with a dynamical variable.

(iii) Probability interpretation of wave function and its implications.

(d) Consider the step potential

$$V(x) = \begin{cases} V_0 & x > 0 \\ 0 & x < 0 \end{cases}$$



Consider a current of particles of mass m propagating from left to right of energy $E > V_0$.

Define $k_1 = \sqrt{\frac{2mE}{\hbar^2}}$, $k_2 = \sqrt{\frac{2m(E - V_0)}{\hbar^2}}$

The general solutions for the regions 1 ($x < 0$) and 2 ($x > 0$) are

$$\phi_1(x) = A_1 e^{ik_1 x} + B_1 e^{-ik_1 x} \quad , \quad \phi_2(x) = A_2 e^{ik_2 x} + B_2 e^{-ik_2 x}$$

(i) State the boundary conditions on the solutions. [2]

(ii) Show that $\frac{B_1}{A_1} = \frac{k_1 - k_2}{k_1 + k_2}$ and $\frac{A_2}{A_1} = \frac{2k_1}{k_1 + k_2}$. [10]

(iii) What is the interpretation of $\frac{B_1}{A_1}$ and $\frac{A_2}{A_1}$? [2]

Q.2. (a) Explain the following:

(i) What is the difference between a state of the system given by the ket $|\ell m\rangle$ and the wave function $\varphi_{\ell m}(r, \vartheta, \varphi)$ describing the same state. [2]

(ii) A dynamical quantity is always represented by a Hermitian linear operator. [2]

(b) (i) A Hamiltonian H has two eigenfunctions ψ_0 and ψ_1 belonging to two different energies E_0 and E_1 respectively. Show that the two eigenfunctions are orthogonal. [5]

(ii) If the two eigenfunctions belong to the same energy, show that they need not be orthogonal. [3]

(c) Consider the one dimensional problem of a particle of mass m in a potential

$$V(x) = \begin{cases} \infty & , \quad x \leq 0, \\ 0 & \quad 0 \leq x \leq a, \\ V_0 & \quad x \geq a. \end{cases}$$

(i) Sketch the potential, [1]

(ii) Find the solutions in the three regions for $E < V_0$, [6]

(iii) Show that the bound state energies ($E < V_0$) are given by the equation [6]

$$\tan\left(\frac{\sqrt{2mE} a}{\hbar}\right) = -\sqrt{\frac{E}{V_0 - E}}$$

Q.3. (a) Verify that the two wave-functions [4]

$$\phi_0(x) = \left(\frac{a}{\pi}\right)^{1/4} \exp(-ax^2/2)$$

$$\phi_1(x) = \left(\frac{a}{\pi}\right)^{1/4} (2a)^{1/2} x \exp(-ax^2/2)$$

are eigenfunctions of the Hamiltonian

$$H = \frac{\hbar^2}{2m} \left(-\frac{d^2}{dx^2} + a^2 x^2 \right)$$

belonging to two different energies.

(i) Show that the two eigenfunctions are normalized. [3]

(ii) Show that $\phi_0(x)$ and $\phi_1(x)$ are orthogonal. [2]

(iii) What are the parities of the two wave functions. [2]

- (iv) Determine the value of $\Delta x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2}$ corresponding to $\phi_0(x)$. [4]

Here $\langle \rangle$ corresponds to expectation value.

- (b) Determine the energies for the ground state $\phi_0(x, y)$ and the first excited state $\phi_1(x, y)$ for the Hamiltonian [10]

$$H = -\frac{\hbar^2}{2m} \left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right] + \frac{m\omega^2}{2} [x^2 + y^2]$$

Use the result

$$-\frac{\hbar^2}{2m} \frac{d^2 u_n(x)}{dx^2} + \frac{m\omega^2}{2} x^2 u_n(x) = (n + \frac{1}{2}) \hbar \omega u_n(x) \text{ where } n=0,1,2,\dots$$

Comment on the degeneracy of the states.

- Q.4. (a)** Using the relation $[x_i, p_j] = i \hbar \delta_{ij}$, where $i, j = x, y, z$ and

$$\vec{L} = \vec{r} \times \vec{p}$$

show that

(i) $[L_z, x] = i\hbar y$; $[L_z, y] = -i\hbar x$ and $[L_z, z] = 0$ [6]

(ii) Using the above identities show that $[L_z, r^2] = 0$ where $r^2 = x^2 + y^2 + z^2$. [2]

- (b) Using the relations

$$[\sigma_x, \sigma_y] = i\hbar \sigma_z, \quad [\sigma_y, \sigma_z] = i\hbar \sigma_x, \quad [\sigma_z, \sigma_x] = i\hbar \sigma_y$$

(i) Show that $[\sigma_i, \sigma^2] = 0$ where $i = x, y, z$ and $\sigma^2 = \sigma_x^2 + \sigma_y^2 + \sigma_z^2$. [6]

- (ii) Consider an Hamiltonian H which has the following properties

$$[H, \sigma^2] = 0$$

$$[H, \sigma_i] = 0$$

Explain why the dynamical variable corresponding to only one component of the operator σ_i can be a measurable simultaneously with H and σ^2 . [3]

- (c) If A and B are two operators which commute with their commutator $[A, B]$, prove that $[A, B^3] = 3B^2 [A, B]$. [8]

Q.5.

- (a) The Hamiltonian of a rotating system with moment of inertia I is given by the expression

$$H = \frac{1}{2I} (L_x^2 + L_y^2)$$

where $\vec{L} = \vec{r} \times \vec{p}$.

- (i) Show that $y_l^m(\vartheta, \varphi)$ are eigenfunctions of H . [2]
(ii) Determine the eigen-value of H . [2]
(iii) Determine the degeneracy of the states for $l = 0, 1, 2$. [6]

(b) An electron is described by an Hamiltonian $H = H_0 + H_1$ where

$$H_0 = \frac{p^2}{2m} - \frac{e^2}{r}$$

and H_1 describes the contribution from an external force.

H_0 has the eigenfunctions $\psi_{nlm} = R_{nl}(r) Y_l^m(\vartheta, \varphi)$ and energy E_n .

Here

n = principal quantum number = 1, 2, 3,

l = orbital angular momentum quantum number = 0, 1, 2, 3,

m = projection of orbital angular momentum or z-component of l .

The eigen functions ψ_{nlm} have following properties:

$$H_0 \psi_{nlm} = E_n \psi_{nlm}$$

$$H_1 \psi_{nlm} = \alpha \psi_{nlm}$$

A state of the electron is described by eigenfunction $\phi = N(\psi_{100} - \sqrt{2} \psi_{210})$

(i) Determine the normalization constant N . [5]

(ii) Determine the expectation value of H . [7]

(iii) Is ϕ an eigenfunction of H ? [3]

$$\text{Note: } \int \psi_{n_1 l_1 m_1}^* \psi_{n_2 l_2 m_2} d\tau = \delta_{n_1 n_2} \delta_{l_1 l_2} \delta_{m_1 m_2}$$

@@@END OF EXAMINATION@@@

APPENDIX:

Given: $h = 6.62606876 \times 10^{-34} \text{ J s}$,

$\hbar = 1.0546 \times 10^{-34} \text{ J s}$,

$c = \text{velocity of light} = 2.99792 \times 10^8 \text{ m s}^{-1}$

$k = 1.3807 \times 10^{-23} \text{ JK}^{-1}$.

$1 \text{ eV} = 1.6022 \times 10^{-19} \text{ J}$,

mass of electron = $9.10938 \times 10^{-31} \text{ kg}$.

mass of neutron/proton = $1.6749 \times 10^{-27} \text{ kg}$,

Useful Information:

$$[A, CD] = [A, C]D + C[A, D]$$

$$[AC, D] = A[C, D] + [A, D]C$$

$$[r_i, p_j] = i \delta_{ij} \text{ where } r_i = (x, y, z) \text{ and } p_i = (p_x, p_y, p_z),$$

$$[L_x, L_y] = i \hbar L_z, [L_y, L_z] = i \hbar L_x, [L_z, L_x] = i \hbar L_y \text{ where } \vec{L} = \vec{r} \times \vec{p},$$

The functions $Y_\ell^m(\vartheta, \varphi)$ are eigenfunctions of L^2 and L_z operators with the property

$$L^2 Y_\ell^m(\vartheta, \varphi) = \ell(\ell+1) \hbar^2 Y_\ell^m(\vartheta, \varphi)$$

$$L_z Y_\ell^m(\vartheta, \varphi) = m \hbar Y_\ell^m(\vartheta, \varphi)$$

Useful Integrals:

$$\int_{-\infty}^{\infty} dz e^{-\alpha z^2} = \sqrt{\frac{\pi}{\alpha}}, \quad \int_{-\infty}^{\infty} dz z^2 e^{-\alpha z^2} = \frac{1}{2} \sqrt{\frac{\pi}{\alpha^3}}, \quad \int_{-\infty}^{\infty} dz z^4 e^{-\alpha z^2} = \frac{3}{4} \sqrt{\frac{\pi}{\alpha^5}}$$

$$\int_0^{\infty} \exp(-t^2) dt = \frac{\sqrt{\pi}}{2}, \quad \int_0^{\infty} t^{2n+1} \exp(-at^2) dt = \frac{n!}{2 a^{n+1}} \quad \text{with } \text{Re } a > 0, n = 0, 1, 2, \dots$$

$$\int_0^{\infty} t^{2n} \exp(-at^2) dt = \frac{1.3.5 \dots (2n-1)}{2^{n+1} a^n} \sqrt{\frac{\pi}{a}} \quad \text{with } \text{Re } a > 0, n = 0, 1, 2, \dots$$

$$\int \sin^2(x) dx = \frac{x}{2} - \frac{1}{4} \sin(2x),$$

$$\int \sin(mx) \sin(nx) dx = \frac{1}{2} \left[\frac{\sin\{(m-n)x\}}{(m-n)} - \frac{\sin\{(m+n)x\}}{(m+n)} \right]$$

$$\int \sin(mx) \cos(nx) dx = -\frac{1}{2} \left[\frac{\cos\{(m-n)x\}}{(m-n)} + \frac{\cos\{(m+n)x\}}{(m+n)} \right]$$

$\int_{-\infty}^{\infty} H_n(\xi) H_m(\xi) \exp(-\xi^2) d\xi = \pi^{\frac{1}{2}} 2^n n! \delta_{nm}$ where $H(\xi)$ are Hermite polynomials and are real.

$$\int_0^{\infty} t^{z-1} \exp(-kt) dt = k^{-z} \Gamma(z) \quad \text{Re } z > 0, \text{Re } k > 0.,$$

$$\Gamma(n+1) = n! \quad \text{for } n = 1, 2, \dots \text{ and } \Gamma(1) = 1.$$

$$\int x^n e^{-m\alpha x} dx = \frac{\Gamma(n+1)}{(m\alpha)^{n+1}} \quad \text{for } m > 0 \text{ and } n \geq 0.$$

You can calculate the integrals you need by expressing powers of x through (repeated) differentiation with respect to the parameter in the exponential, e.g.

$$\int_a^b dx x \exp(-\gamma x) = -\frac{\partial}{\partial \gamma} \int_a^b dx \exp(-\gamma x) \quad \text{and} \quad \int_a^b dx x^2 \exp(-\gamma x) = \frac{\partial^2}{\partial \gamma^2} \int_a^b dx \exp(-\gamma x) \quad \text{and so on.}$$