

**UNIVERSITY OF SWAZILAND**

**FACULTY OF SCIENCE AND ENGINEERING**

**DEPARTMENT OF PHYSICS**

**MAIN EXAMINATION 2012/2013**

**TITLE OF PAPER : CLASSICAL MECHANICS**

**COURSE NUMBER : P320**

**TIME ALLOWED : THREE HOURS**

**INSTRUCTIONS : ANSWER ANY FOUR OUT OF FIVE  
QUESTIONS.  
EACH QUESTION CARRIES 25  
MARKS.  
MARKS FOR DIFFERENT SECTIONS  
ARE SHOWN IN THE RIGHT-HAND  
MARGIN.**

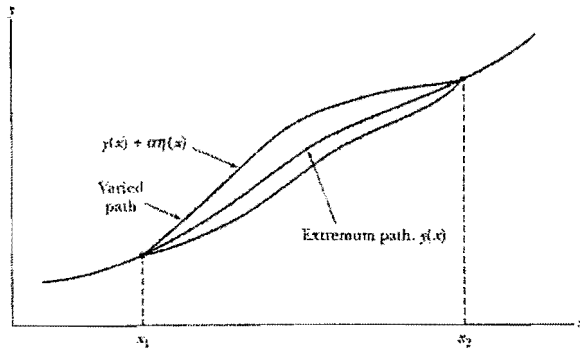
**THIS PAPER HAS EIGHT PAGES, INCLUDING THIS PAGE.**

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**P320 CLASSICAL MECHANICS**

**Question one**

- (a) Given the following definite integral  $J(\alpha) = \int_{x_1}^{x_2} f(y(\alpha, x), y'(\alpha, x); x) dx$ , where the varied integration path is  $y(\alpha, x) = y(x) + \alpha \eta(x)$  and  $\eta(x_1) = \eta(x_2) = 0$  as shown in the following diagram :



Using the extremum condition for  $J(\alpha)$ , i.e.,  $\left. \frac{\partial J(\alpha)}{\partial \alpha} \right|_{\alpha=0} = 0$ , to deduce that

$f$  along the extremum path, i.e.,  $f(y(x), y'(x); x)$ , satisfies the following equation:

$$\frac{\partial f}{\partial y} - \frac{d}{dx} \left( \frac{\partial f}{\partial y'} \right) = 0 \quad (10 \text{ marks})$$

- (b) If  $H$  denotes the Hamiltonian function and  $L$  is the Lagrangian function, use the definition  $H = \sum_{\alpha=1}^n p_{\alpha} \dot{q}_{\alpha} - L$  (where  $p_{\alpha}$  and  $q_{\alpha}$  ( $\alpha = 1, 2, \dots, n$ ) are the generalized momenta and coordinates respectively, i.e.,  $H = H(q_1, \dots, q_n, p_1, \dots, p_n, t)$ ,

$L = L(q_1, \dots, q_n, \dot{q}_1, \dots, \dot{q}_n, t)$ ,  $p_{\alpha} = \frac{\partial L}{\partial \dot{q}_{\alpha}}$  and  $\dot{p}_{\alpha} = \frac{\partial L}{\partial q_{\alpha}}$ ) to show that

(i)  $\dot{q}_{\alpha} = \frac{\partial H}{\partial p_{\alpha}} \quad \alpha = 1, 2, \dots, n \quad (4 \text{ marks})$

(ii)  $\dot{p}_{\alpha} = -\frac{\partial H}{\partial q_{\alpha}} \quad \alpha = 1, 2, \dots, n \quad (4 \text{ marks})$

(iii)  $\frac{\partial H}{\partial t} = -\frac{\partial L}{\partial t} \quad (7 \text{ marks})$

### Question two

- (a) The Poisson Bracket  $[F, G]$  of two functions  $F$  and  $G$  of canonical variables  $p_\alpha$  and  $q_\alpha$  is given by

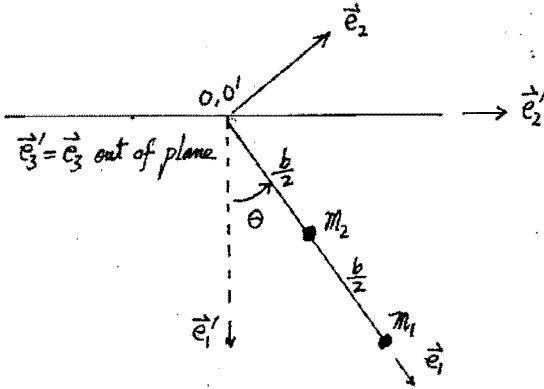
$$[F, G] \equiv \sum_{\alpha=1}^n \left( \frac{\partial F}{\partial q_\alpha} \frac{\partial G}{\partial p_\alpha} - \frac{\partial F}{\partial p_\alpha} \frac{\partial G}{\partial q_\alpha} \right)$$

Show that

(i)  $[H, q_\alpha] = -\dot{q}_\alpha$  where  $H$  is the Hamiltonian, ( 4 marks )

(ii)  $[p_\alpha, p_\beta] = 0$  where  $\begin{cases} \alpha = 1, 2, \dots, n \\ \beta = 1, 2, \dots, n \end{cases}$ , ( 3 marks )

- (b) A pendulum is composed of a rigid rod of length  $b$  with a mass  $m_1$  at its end. Another mass  $m_2$  is placed halfway down the rod. The mass of the rod itself is negligible. Let the fixed and the body coordinate systems have their origin at the pendulum pivot point. Let  $(\vec{e}_1', \vec{e}_2', \vec{e}_3')$  and  $(\vec{e}_1, \vec{e}_2, \vec{e}_3)$  be the unit vectors of the fixed and the body coordinate system respectively as shown below.



- (i) Write down the inertia tensor  $I$  for the pendulum with respect to the body coordinate system given above and deduce that  $I$  is a diagonal matrix with its diagonal elements as  $I_{1,1} = 0$  and  $I_{2,2} = I_{3,3} = \left( m_1 + \frac{m_2}{4} \right) b^2$ . ( 6 marks )

- (ii) From the equation of rotational motion, i.e.,  $\dot{\vec{L}} = \vec{N}$  where angular momentum  $\vec{L} = I \vec{\omega}$  and torque  $\vec{N} = \sum_a (\vec{r}_a \times \vec{F}_a)$ , deduce the following equation :

$$b^2 \left( m_1 + \frac{m_2}{4} \right) \ddot{\theta} = - b g \sin(\theta) \left( m_1 + \frac{m_2}{2} \right) \quad ( 12 \text{ marks } )$$

(Hint :

$$\vec{\omega} = \vec{e}_3' \dot{\theta} \quad , \quad \vec{F}_1 = \vec{e}_1' m_1 g \quad , \quad \vec{F}_2 = \vec{e}_1' m_2 g \quad ,$$

$$\vec{r}_1 = \vec{e}_1' b \cos(\theta) + \vec{e}_2' b \sin(\theta) \quad \text{and} \quad \vec{r}_2 = \vec{e}_1' \frac{b}{2} \cos(\theta) + \vec{e}_2' \frac{b}{2} \sin(\theta) \quad )$$

### Question three

- (a) For circular orbits in an attractive central force potential of the form  $V = -\frac{k}{r^n}$  where  $k$  is a positive constant and  $n > 0$ , find a relation between the kinetic and potential energies and show that

$$T = \frac{n k}{2 r^n} \quad (9 \text{ marks})$$

(Hint :  $\vec{a} = \vec{e}_r (\ddot{r} - r \dot{\theta}^2) + \vec{e}_\theta (2 \dot{r} \dot{\theta} + r \ddot{\theta})$ )

- (b) Starting from the law of conservation of angular momentum  $l$ , derive Kepler's third law, i.e., the relation between the period  $\tau$  of a closed orbit in an attractive inverse square central force and the area  $A$  of the orbit. Show that

$$\tau = \frac{2 \mu}{l} A \quad \text{where } \mu \text{ is the reduced mass of the system.} \quad (7 \text{ marks})$$

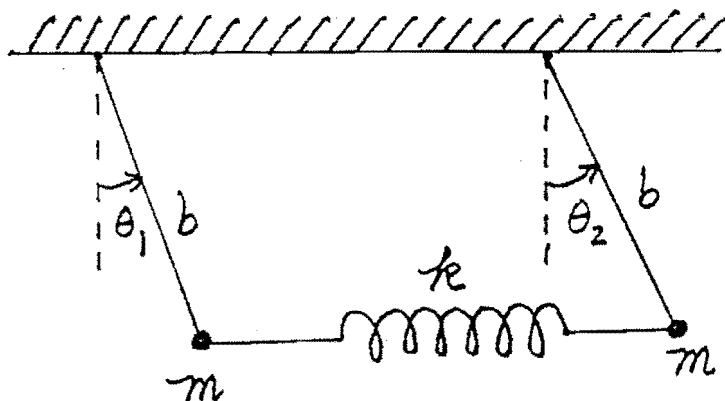
- (c) An earth satellite moves in an elliptical orbit with period  $\tau$ , eccentricity  $\varepsilon$  and semi-major axis  $a$ . The maximum radial velocity, named as  $v_{\theta, \max}$ , occurs at  $r = r_{\min}$ . Show that

$$v_{\theta, \max} = \frac{2 \pi a}{\tau \sqrt{1 - \varepsilon^2}} \quad (9 \text{ marks})$$

(Hint :  $\mu r_{\min} v_{\theta, \max} = l$ ,  $A = \pi a b$  and  $b = a \sqrt{1 - \varepsilon^2}$ )

### Question four

Two pendulums of equal lengths  $b$  and equal masses  $m$  are connected by a spring of force constant  $k$  as shown below. The spring is unstretched in the equilibrium position, i.e.,  $\theta_1 = 0$  and  $\theta_2 = 0$ .



- (i) For small  $\theta_1$  and  $\theta_2$ , i.e.,

$\left( \sin(\theta_1) \approx \theta_1, \sin(\theta_2) \approx \theta_2, \cos(\theta_1) \approx 1 - \frac{\theta_1^2}{2} \text{ and } \cos(\theta_2) \approx 1 - \frac{\theta_2^2}{2} \right)$ , show that the

Lagrangian for the system can be expressed as:

$$L = \frac{1}{2} m b^2 (\dot{\theta}_1^2 + \dot{\theta}_2^2) - \frac{m g b}{2} (\theta_1^2 + \theta_2^2) - \frac{k b^2}{2} (\theta_1 - \theta_2)^2$$

where the zero gravitational potential is set at the equilibrium position. **(8 marks)**

- (ii) Write down the equations of motion and deduce that

$$\begin{cases} \ddot{\theta}_1 = -\left(\frac{m g + k b}{m b}\right) \theta_1 + \frac{k}{m} \theta_2 \\ \ddot{\theta}_2 = \frac{k}{m} \theta_1 - \left(\frac{m g + k b}{m b}\right) \theta_2 \end{cases} \quad \text{(6 marks)}$$

- (iii) Set  $\theta_1 = \hat{X}_1 e^{i\omega t}$  and  $\theta_2 = \hat{X}_2 e^{i\omega t}$  (where  $\hat{X}_1$  and  $\hat{X}_2$  are constants) and deduce from the equations in (ii) the matrix equation  $-\omega^2 X = A X$  where

$$X = \begin{pmatrix} \hat{X}_1 \\ \hat{X}_2 \end{pmatrix} \text{ and } A = \begin{pmatrix} -\left(\frac{m g + k b}{m}\right) & \frac{k b}{m} \\ \frac{k b}{m} & -\left(\frac{m g + k b}{m}\right) \end{pmatrix} \quad \text{(4 marks)}$$

- (iv) Find the eigenfrequencies  $\omega$  of this coupled system and show that they are

$$\omega = \sqrt{\frac{g}{b}} \quad \text{or} \quad \sqrt{\frac{m g + 2 k b}{m b}} \quad \text{(7 marks)}$$

### Question five

- (a) Two sets of coordinate systems are having the same origins. The non-prime system (with position vector denoted as  $\vec{r}$  and referred to as “rotating” system) is rotating with an angular velocity  $\vec{\omega}$  about the prime system (with position vector denoted as  $\vec{r}'$  and referred as “fixed” system and taken as an inertial system). Use the following proven

relation that  $\left(\frac{d\vec{F}}{dt}\right)_{fixed} = \left(\frac{d\vec{F}}{dt}\right)_{rotating} + \vec{\omega} \times \vec{F}$  for any vector field  $\vec{F}$  to deduce the

following:

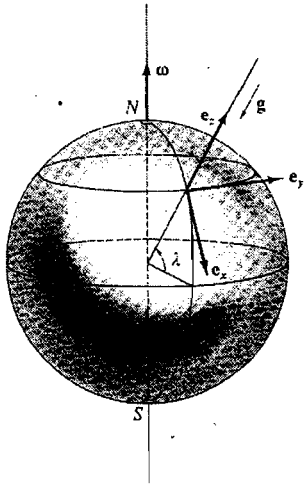
$$\vec{a}_{eff} = \vec{a} - \dot{\vec{\omega}} \times \vec{r} - \vec{\omega} \times (\vec{\omega} \times \vec{r}) - 2\vec{\omega} \times \vec{v}_r \quad \text{where}$$

$$\vec{r}' = \vec{r} \text{ (same origin)} \quad , \quad \vec{a}_{eff} \equiv \left(\frac{d^2\vec{r}}{dt^2}\right)_{rotating} \quad , \quad \vec{a} \equiv \left(\frac{d^2\vec{r}'}{dt^2}\right)_{fixed} \quad , \quad \vec{v}_r \equiv \left(\frac{d\vec{r}}{dt}\right)_{rotating} \quad ,$$

$$\left(\frac{d\vec{v}_r}{dt}\right)_{rotating} \equiv \left(\frac{d^2\vec{r}}{dt^2}\right)_{rotating} \quad \text{and} \quad \dot{\vec{\omega}} \equiv \left(\frac{d\vec{\omega}}{dt}\right)_{rotating}$$

( 12 marks )

- (b)



Show that the horizontal deflection  $d$  along  $-\vec{e}_y$  direction resulting from the Coriolis force  $(-2m\vec{\omega} \times \vec{v}_r)$  of a particle falling freely in the earth's gravitational field at a northern latitude  $\lambda$  is

$$d \approx \frac{1}{3} \omega \cos(\lambda) \sqrt{\frac{8h^3}{g}} \quad \text{where}$$

$\omega$  : angular velocity of earth's rotation

( 13 marks )

$h$  : the height of the particle above the earth before its free fall

(Hint :

$$\vec{a}_{eff} \approx \vec{e}_z (-g) - 2\vec{\omega} \times \vec{v}_r \quad , \quad \vec{v}_r \approx \vec{e}_z (-gt) \quad , \quad \vec{\omega} = \vec{e}_x (-\omega \cos(\lambda)) + \vec{e}_z (\omega \sin(\lambda))$$

$$\text{and (total time for the given motion)} = \sqrt{\frac{2h}{g}} \quad )$$

Useful informations

$$V = - \int \vec{F} \bullet d\vec{l} \quad \text{and reversely} \quad \vec{F} = -\vec{\nabla} V$$

$$L = T - V = L(q_1, q_2, \dots, q_n, \dot{q}_1, \dot{q}_2, \dots, \dot{q}_n, t)$$

$$p_\alpha = \frac{\partial L}{\partial \dot{q}_\alpha} \quad \text{and} \quad \dot{p}_\alpha = \frac{\partial L}{\partial q_\alpha}$$

$$H = \sum_{\alpha=1}^n (p_\alpha \dot{q}_\alpha) - L = H(q_1, q_2, \dots, q_n, \dot{q}_1, \dot{q}_2, \dots, \dot{q}_n, t)$$

$$\dot{q}_\alpha = \frac{\partial H}{\partial p_\alpha} \quad \text{and} \quad \dot{p}_\alpha = -\frac{\partial H}{\partial q_\alpha}$$

$$[u, v] \equiv \sum_{\alpha=1}^n \left( \frac{\partial u}{\partial q_\alpha} \frac{\partial v}{\partial p_\alpha} - \frac{\partial u}{\partial p_\alpha} \frac{\partial v}{\partial q_\alpha} \right)$$

$$G = 6.673 \times 10^{-11} \frac{N m^2}{kg^2}$$

$$\text{radius of earth } r_E = 6.4 \times 10^6 \text{ m}$$

$$\text{mass of earth } m_E = 6 \times 10^{24} \text{ kg}$$

$$\text{earth attractive potential} \equiv -\frac{k}{r} \quad \text{where} \quad k = G m m_E$$

$$\varepsilon = \sqrt{1 + \frac{2 E l^2}{\mu k}} \quad \{(\varepsilon = 0, \text{circle}), (0 < \varepsilon < 1, \text{ellipse}), (\varepsilon = 1, \text{parabola}), \dots\}$$

$$\mu = \frac{m_1 m_2}{m_1 + m_2} \approx m_1 \quad \text{if} \quad m_2 \gg m_1$$

$$\text{For elliptical orbit, i.e., } 0 < \varepsilon < 1, \text{ then} \left\{ \begin{array}{l} \text{semi-major } a = \frac{k}{2|E|} \\ \text{semi-minor } b = \frac{l}{\sqrt{2\mu|E|}} \\ \text{period } \tau = \frac{2\mu}{l} (\pi a b) \\ r_{\min} = a(1 - \varepsilon) \quad \& \quad r_{\max} = a(1 + \varepsilon) \end{array} \right.$$

for plane polar  $(r, \theta)$  system with unit vectors  $(\vec{e}_r, \vec{e}_\theta)$ , we have

$$\left\{ \begin{array}{l} \vec{v} = \vec{e}_r \dot{r} + \vec{e}_\theta r \dot{\theta} \\ \vec{a} = \vec{e}_r (\ddot{r} - r \dot{\theta}^2) + \vec{e}_\theta (2\dot{r} \dot{\theta} + r \ddot{\theta}) \end{array} \right.$$

$$\vec{\nabla} f = \vec{e}_r \frac{\partial f}{\partial r} + \vec{e}_\theta \frac{1}{r} \frac{\partial f}{\partial \theta}$$

**Useful informations (continued)**

$$I = \begin{pmatrix} \sum_{\alpha} m_{\alpha} (x_{\alpha,2}^2 + x_{\alpha,3}^2) & -\sum_{\alpha} m_{\alpha} x_{\alpha,1} x_{\alpha,2} & -\sum_{\alpha} m_{\alpha} x_{\alpha,1} x_{\alpha,3} \\ -\sum_{\alpha} m_{\alpha} x_{\alpha,2} x_{\alpha,1} & \sum_{\alpha} m_{\alpha} (x_{\alpha,1}^2 + x_{\alpha,3}^2) & -\sum_{\alpha} m_{\alpha} x_{\alpha,2} x_{\alpha,3} \\ -\sum_{\alpha} m_{\alpha} x_{\alpha,3} x_{\alpha,1} & -\sum_{\alpha} m_{\alpha} x_{\alpha,3} x_{\alpha,2} & \sum_{\alpha} m_{\alpha} (x_{\alpha,1}^2 + x_{\alpha,2}^2) \end{pmatrix}$$

$$\vec{F}_{eff} = \vec{F} - m \ddot{\vec{R}}_f - m \dot{\vec{\omega}} \times \vec{r} - m \vec{\omega} \times (\vec{\omega} \times \vec{r}) - 2 m \vec{\omega} \times \vec{v}_r, \quad \text{where}$$

$$\vec{r}' = \vec{R} + \vec{r} \quad \text{and}$$

$\vec{r}'$  refers to fixed (inertial system)

$\vec{r}$  refers to rotational (non-inertial system) rotates with  $\vec{\omega}$  to  $\vec{r}'$  system

$\vec{R}$  from the origin of  $\vec{r}'$  to the origin of  $\vec{r}$

$$\vec{v}_r = \left( \frac{d\vec{r}}{dt} \right)_r$$