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UNIVERSITY OF SWAZILAND
FACULTY OF SCIENCE
DEPARTMENT OF PHYSICS
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## TITLE OF PAPER : CLASSICAL MECHANICS

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COURSE NUMBER : P320
TIME ALLOWED : THREE HOURS
INSTRUCTIONS : ANSWER ANY FOUR OUT OF FIVE QUESTIONS.
EACH QUESTION CARRIES 25 MARKS.
MARKS FOR DIFFERENT SECTIONS ARE SHOWN IN THE RIGHT-HAND MARGIN.
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## Question one

(a) If $H$ denotes the Hamiltonian function and $L$ is the Lagrangian function, use the definition $H=\sum_{\alpha=1}^{n} p_{\alpha} \dot{q}_{\alpha}-L \quad$ [where $p_{\alpha}$ and $q_{\alpha}(\alpha=1,2, \cdots, n)$ are the generalized momenta and coordinates respectively, ie., $H=H\left(q_{1}, \cdots, q_{n}, p_{1}, \cdots, p_{n}, t\right)$, $L=L\left(q_{1}, \cdots, q_{n}, \dot{q}_{1}, \cdots, \dot{q}_{n}, t\right) \quad, \quad p_{\alpha}=\frac{\partial L}{\partial \dot{q}_{\alpha}}$ and $\left.\dot{p}_{\alpha}=\frac{\partial L}{\partial q_{\alpha}}\right]$ to show that
(i) $\quad \dot{q}_{\alpha}=\frac{\partial H}{\partial p_{\alpha}} \quad \alpha=1,2, \cdots, n$
(4 marks)
(ii) $\quad \dot{p}_{\alpha}=-\frac{\partial H}{\partial q_{\alpha}} \quad \alpha=1,2, \cdots, n$
(iii) $\frac{\partial H}{\partial t}=-\frac{\partial L}{\partial t}$
(4 marks)
( 7 marks )
(b) For a certain dynamical system the kinetic energy, $T$, and potential energy, $V$, are given by $T=\dot{q}_{1}^{2}+2 \dot{q}_{1} \dot{q}_{2}+3 \dot{q}_{2}^{2}$ $V=4 q_{1}^{2}$
where $q_{1}, q_{2}$ are the generalized coordinates.
(i) Find the momentum $p_{1} \& p_{2}$ of the system.
(ii) Use $H=\sum_{\alpha=1}^{2} p_{\alpha} \dot{q}_{\alpha}-L$ to find the Hamiltonian function of the system and show that $H=\frac{1}{8}\left(3 p_{1}{ }^{2}-2 p_{1} p_{2}+p_{2}{ }^{2}\right)+4 q_{1}^{2}$

## Question two

The definition of the Poisson brackets are given as $[u, v]_{q, p} \equiv \sum_{\alpha=1}^{n}\left(\frac{\partial u}{\partial q_{\alpha}} \frac{\partial v}{\partial p_{\alpha}}-\frac{\partial u}{\partial p_{\alpha}} \frac{\partial v}{\partial q_{\alpha}}\right)$, or simply written as $[u, v]$, where $q_{\alpha}$ and $p_{\alpha}$ are the $\alpha^{i h}$ generalized coordinate and momentum respectively .
(a) For any function $F\left(q_{1}, q_{2}, \cdots, q_{n}, p_{1}, p_{2}, \cdots, p_{n}, t\right)$, prove that $\frac{d F}{d t}=[F, H]+\frac{\partial F}{\partial t}$
where $H$ is the Hamiltonian of the system, i.e., $H\left(q_{1}, q_{2}, \cdots, q_{n}, p_{1}, p_{2}, \cdots, p_{n}, t\right)$
( 5 marks)
(b) The three components of the angular momentum $\vec{l}(\equiv \vec{r} \times \vec{p})$ of a particle of mass $m$ are given by $l_{1}=q_{2} p_{3}-q_{3} p_{2}, l_{2}=q_{3} p_{1}-q_{1} p_{3}$ and $l_{3}=q_{1} p_{2}-q_{2} p_{1}$ where $\quad p_{i}=m \dot{q}_{i} \quad i=1,2,3$. Show that
(i) $\left[l_{1}, l_{2}\right]=l_{3}$
( 5 marks)
(ii) $\left[q_{2}, l_{3}\right]=q_{1}$
( 5 marks)
(c) For an equation of the type $\frac{d u}{d t}=[u, H]$ the specific solution of $u(t)$ is given by the following Taylor series expansion for the time $t$ as

$$
\begin{equation*}
\left.\left.u(t)=u_{0}+[u, H]_{0} t+[\llbracket u, H], H\right]_{0} \frac{t^{2}}{2!}+[[u, H], H], H\right]_{0} \frac{t^{3}}{3!}+\cdots \cdots \cdots \tag{1}
\end{equation*}
$$

where subscript 0 denotes the initial conditions at $t=0$.
For a simple harmonic oscillator system described by $H=\frac{p^{2}}{2 m}+\frac{1}{2} k x^{2}$, if the initial conditions are given as $x_{0}$ and $p_{0}$, use eq.(1) to deduce that

$$
\begin{equation*}
x(t)=x_{0}+\frac{p_{0}}{m} t-\frac{k x_{0}}{2 m} t^{2}-\frac{k p_{0}}{6 m^{2}} t^{3}+\cdots \cdots \tag{10marks}
\end{equation*}
$$

## Question three

(a) Given the Lagrangian for the two-body central force system as:
$L=T-V=\frac{1}{2} \mu\left(\dot{r}^{2}+r^{2} \dot{\theta}^{2}\right)+\frac{k}{r}$
where $\mu$ is the reduced mass of the system, $k$ is a positive constant and $(r, \theta)$ are polar coordinates of the motion plane with its origin at the center of mass of the two-body system.
(i) Write down the Lagrange's equation for $\theta$ and show that the angular momentum $l$ is conserved, ie., deduce that
$\dot{\theta}=\frac{l}{\mu r^{2}} \quad \cdots \cdots$ (1) where $l$ is a constant.
( 3 marks )
(ii) Write down the Lagrange's equation for $r$, with eq.(1) inserted, deduce that $\mu \ddot{r}-\frac{l^{2}}{\mu r^{3}}+\frac{k}{r^{2}}=0$
( 3 marks )
(iii) Multiply eq.(2) by $d r$ and use $\ddot{r} d r=\frac{d \dot{r}}{d t} d r=d \dot{r} \frac{d r}{d t}=\dot{r} d \dot{r}=d\left(\frac{\dot{r}^{2}}{2}\right)$ to show that the total energy $E(\equiv T+V)$ is conserved, i.e., $\frac{1}{2} \mu\left(\dot{r}^{2}+r^{2} \dot{\theta}^{2}\right)-\frac{k}{r}=$ const.$\equiv E$
( 6 marks )
(b) If an earth satellite of 500 kg mass is having a pure tangential speed $v_{\theta}=8,000 \mathrm{~m} / \mathrm{s}$ at its near-earth-point 600 km above the earth surface,
(i) calculate the values of $l$ and $E$ of this satellite, ( 3 marks)
(ii) calculate the values of the eccentricity, $\varepsilon$, and show that the orbit is an elliptical orbit. Also calculate its period.
( 6 marks)
(iii) determine the value of the $v_{\theta}$ at the same given near-earth-point such that the satellite orbit is a circular orbit,
( 2 marks)
(Hint : $E=\frac{1}{2} \mu v_{\theta}^{2}-\frac{k}{r} \xrightarrow{\text { circular orbit }}-\frac{k}{2 r}$ )
(iv) determine the value of the $v_{\theta}$ at the same given near-earth-point such that the satellite orbit is a parabolic orbit .
(a) Two set of Cartesian coordinate axes are having the same origins and $z$-axis. The nonprime system (referred to as "rotating" system) is rotating with an angular velocity $\vec{\omega}=\vec{e}_{z^{\prime}} \dot{\theta}$ about the prime system (referred as "fixed" system) as shown below:


For any vector field $\vec{F}$ decomposed into the above two-set of cartesian components, i.e., $\vec{F}=\vec{e}_{x} F_{x}+\vec{e}_{y} F_{y}+\vec{e}_{z} F_{z}=\vec{e}_{x^{\prime}} F_{x^{\prime}}+\vec{e}_{y^{\prime}} F_{y^{\prime}}+\vec{e}_{z^{\prime}} F_{z^{\prime}}$, show that $\left(\frac{d \vec{F}}{d t}\right)_{\text {fixed }}=\left(\frac{d \vec{F}}{d t}\right)_{\text {rotuting }}+\vec{a} \times \vec{F} \quad$ where
$\left(\frac{d \vec{F}}{d t}\right)_{\text {fured }}=\vec{e}_{x^{\prime}} \frac{d F_{x^{\prime}}}{d t}+\vec{e}_{y^{\prime}} \frac{d F_{y^{\prime}}}{d t}+\vec{e}_{z^{\prime}} \frac{d F_{z^{\prime}}}{d t}$ and
$\left(\frac{d \vec{F}}{d t}\right)_{\text {rotating }}=\vec{e}_{x} \frac{d F_{x}}{d t}+\vec{e}_{y} \frac{d F_{y}}{d t}+\vec{e}_{z} \frac{d F_{z}}{d t}$
(Hint : $\vec{e}_{x}=\vec{e}_{x^{\prime}} \cos (\theta)+\vec{e}_{y^{\prime}} \sin (\theta), \vec{e}_{y}=-\vec{e}_{x^{\prime}} \sin (\theta)+\vec{e}_{y^{\prime}} \cos (\theta)$ and $\vec{e}_{x}=\vec{e}_{z^{\prime}}$ )
(b)


If a person, near the earth surface at a northern latitude $\lambda$, fired a bullet of speed, $v_{0}$, at a target situated at his north direction ( $-\vec{e}_{x}$ direction) of distance $L$ away from him. Assuming he has a perfect rifle and the time $T$ for the bullet hitting the target is short and $T \approx \frac{L}{v_{0}}$ (i.e., neglecting the gravitational bending and assuming the bullet is moving along - $x$ direction with constant speed $v_{0}$ ).
(i) Show that the bullet will miss the target by a deviation distance $d$ resulting from the Coriolis force $\left(-2 m \vec{\omega} \times \vec{v}_{r}\right)$. Show that

$$
d=\frac{\omega L^{2}}{v_{0}} \sin (\lambda)
$$

( 11 marks )
(Hint: $\left.\vec{a}_{e f f} \approx-2 \vec{\omega} \times \vec{v}_{r}, \vec{v}_{r} \approx \vec{e}_{x}\left(-v_{0}\right), \vec{\omega}=\vec{e}_{x}(-\omega \cos (\lambda))+\vec{e}_{z}(\omega \sin (\lambda))\right)$
(ii) Given the values of $\lambda=60^{\circ}, L=2000 \mathrm{~m}, \nu_{0}=800 \mathrm{~m} / \mathrm{s}$ and $\omega=2 \pi \mathrm{rad} / \mathrm{day}$ (i.e., $\omega=7.27 \times 10^{-5} \mathrm{rad} / \mathrm{s}$ ), determine the value of the deviation distance $d$
( 2 marks)

## Question five

Six equal mass point $m\left(=m_{1}=m_{2}=\cdots=m_{6}\right)$ attached by massless rigid rods to form a rigid body of diamond shape with the center of mass of the "diamond" chosen as the origin of the body coordinate system $\left(x_{1}, x_{2}, x_{3}\right)$ as shown in the diagram below.

where each mass point's coordinates in terms of length $a \& b$ is indicated in the diagram.
(a) Evaluate all elements of the inertia tensor, $I$, of the given rigid body with respect to the chosen body coordinate system and show that

$$
I=\left(\begin{array}{ccc}
2 m\left(a^{2}+b^{2}\right) & 0 & 0 \\
0 & 2 m\left(a^{2}+b^{2}\right) & 0 \\
0 & 0 & 4 m a^{2}
\end{array}\right)
$$

( 6 marks )
(b) If the given rigid body is only rotating with an angular velocity $\vec{\omega}$ without translational motion with respect to a fixed inertia coordinate system $\left(x_{1}^{\prime}, x_{2}^{\prime}, x_{3}^{\prime}\right)$ sharing the same origin as that of the body coordinate system, write down the total kinetic energy
$T=T_{\text {roational }}=\frac{1}{2} \vec{\omega} \bullet I \bullet \vec{\omega}$ in terms of $m, a, b, \omega_{1}, \omega_{2} \& \omega_{3} \quad$ where

$$
\begin{equation*}
\vec{\omega}=\vec{e}_{1} \omega_{1}+\vec{e}_{2} \omega_{2}+\vec{e}_{3} \omega_{3} \tag{2marks}
\end{equation*}
$$

(c) The following are Euler's equations for force-free pure-rotational motion, i.e., $L=T_{\text {rotational }}$, for already diagonalized $I$ as the case in (a).
$\left\{\begin{array}{llll}\left(I_{2}-I_{3}\right) \omega_{2} \omega_{3}-I_{1} \dot{\omega}_{1}=0 & \cdots \cdots & \text { (1) } \\ \left(I_{3}-I_{1}\right) \omega_{3} \omega_{1}-I_{2} \dot{\omega}_{2}=0 & \cdots \cdots & \text { (2) } \\ \left(I_{1}-I_{2}\right) \omega_{1} \omega_{2}-I_{3} \dot{\omega}_{3}=0 & \cdots \cdots & \text { (3) }\end{array}\right.$
(i) For our given rigid body, deduce from the above Euler's equations that

$$
\left\{\begin{array}{c}
\omega_{3}=\text { const. } \xrightarrow{\text { setas }} K  \tag{4}\\
\dot{\omega}_{1}=\frac{\left(-a^{2}+b^{2}\right) K}{\left(a^{2}+b^{2}\right)} \omega_{2} \\
\dot{\omega}_{2}=-\frac{\left(-a^{2}+b^{2}\right) K}{\left(a^{2}+b^{2}\right)} \omega_{1}
\end{array}\right.
$$

(ii) If $b>a \& K>0$, then $\frac{\left(-a^{2}+b^{2}\right) K}{\left(a^{2}+b^{2}\right)}$ is a positive constant and set it as $\Omega$. Deduce from eq.(5) and eq.(6) in (c)(i) that $\ddot{\omega}_{1}=-\Omega^{2} \omega_{1}$
( 3 marks)
(iii) By direct substitution, show that $\omega_{1}=A \cos (\Omega t+B) \cdots \cdots$ (8) is the solution to eq.(7) with $A \& B$ constant values linking to the given initial value of $\bar{\omega}$.
( 2 marks)
(iv) Substitute eq.(8) into eq.(5) and deduce that

$$
\begin{equation*}
\omega_{2}=-A \sin (\Omega t+B) \tag{9}
\end{equation*}
$$

(v) Show that the magnitude of $\vec{\omega}$ is a constant for all time $t$.

## Useful informations

$V=-\int \vec{F} \cdot d \vec{l}$ and reversely $\vec{F}=-\vec{\nabla} V$
$L=T-V=L\left(q_{1}, q_{2}, \cdots, q_{n}, \dot{q}_{1}, \dot{q}_{2}, \cdots, \dot{q}_{n}, t\right)$
$p_{\alpha}=\frac{\partial L}{\partial \dot{q}_{\alpha}} \quad$ and $\quad \dot{p}_{\alpha}=\frac{\partial L}{\partial q_{\alpha}}$
$H=\sum_{\alpha=1}^{n}\left(p_{\alpha} \dot{q}_{\alpha}\right)-L=H\left(q_{1}, q_{2}, \cdots, q_{n}, \dot{q}_{1}, \dot{q}_{2}, \cdots, \dot{q}_{n}, t\right)$
$\dot{q}_{\alpha}=\frac{\partial H}{\partial p_{\alpha}} \quad$ and $\quad \dot{p}_{\alpha}=-\frac{\partial H}{\partial q_{\alpha}}$
$[u, v] \equiv \sum_{\alpha=1}^{n}\left(\frac{\partial u}{\partial q_{\alpha}} \frac{\partial v}{\partial p_{\alpha}}-\frac{\partial u}{\partial p_{\alpha}} \frac{\partial v}{\partial q_{\alpha}}\right)$
$G=6.673 \times 10^{-11} \frac{\mathrm{Nm}^{2}}{\mathrm{~kg}^{2}}$
radius of earth $r_{E}=6.4 \times 10^{6} \mathrm{~m}$
mass of earth $m_{E}=6 \times 10^{24} \mathrm{~kg}$
earth attractive potential $\equiv-\frac{k}{r} \quad$ where $\quad k=G m m_{E}$
$\varepsilon=\sqrt{1+\frac{2 E l^{2}}{\mu k}}\{(\varepsilon=0$, circle $),(0<\varepsilon<1$, ellipse $),(\varepsilon=1$, parabola $), \cdots\}$
$\mu=\frac{m_{1} m_{2}}{m_{1}+m_{2}} \approx m_{1} \quad$ if $\quad m_{2} \gg m_{1}$
For elliptical orbit,i.e., $0<\varepsilon<1$, then $\left\{\begin{array}{c}\text { semi-major } a=\frac{k}{2|E|} \\ \text { semi-minor } b=\frac{l}{\sqrt{2 \mu|E|}} \\ \text { period } \tau=\frac{2 \mu}{l}(\pi a b) \\ r_{\min }=a(1-\varepsilon) \& r_{\max }=a(1+\varepsilon)\end{array}\right.$
for plane polar $(r, \theta)$ system with unit vectors $\left(\vec{e}_{r}, \vec{e}_{\theta}\right)$, we have
$\left\{\begin{array}{l}\vec{v}=\vec{e}_{r} \dot{r}+\vec{e}_{\theta} r \dot{\theta} \\ \vec{a}=\vec{e}_{r}\left(\ddot{r}-r \dot{\theta}^{2}\right)+\vec{e}_{\theta}(2 \dot{r} \dot{\theta}+r \ddot{\theta})\end{array}\right.$
$\vec{\nabla} f=\vec{e}_{r} \frac{\partial f}{\partial r}+\vec{e}_{\theta} \frac{1}{r} \frac{\partial f}{\partial \theta}$

## Useful informations (continued)

$I=\left(\begin{array}{ccc}\sum_{\alpha} m_{\alpha}\left(x_{\alpha, 2}^{2}+x_{\alpha, 3}^{2}\right) & -\sum_{\alpha} m_{\alpha} x_{\alpha, 1} x_{\alpha, 2} & -\sum_{\alpha} m_{\alpha} x_{\alpha, 1} x_{\alpha, 3} \\ -\sum_{\alpha} m_{\alpha} x_{\alpha, 2} x_{\alpha, 1} & \sum_{\alpha} m_{\alpha}\left(x_{\alpha, 1}^{2}+x_{\alpha, 3}^{2}\right) & -\sum_{\alpha} m_{\alpha} x_{\alpha, 2} x_{\alpha, 3} \\ -\sum_{\alpha}^{2} m_{\alpha} x_{\alpha, 3} x_{\alpha, 1} & -\sum_{\alpha} m_{\alpha} x_{\alpha, 3} x_{\alpha, 2} & \sum_{\alpha} m_{\alpha}\left(x_{\alpha, 1}^{2}+x_{\alpha, 2}^{2}\right)\end{array}\right)$
$\vec{F}_{e f f}=\vec{F}-m \ddot{\vec{R}}_{f}-m \dot{\vec{\omega}} \times \vec{r}-m \vec{\omega} \times(\vec{\omega} \times \vec{r})-2 m \vec{\omega} \times \vec{v}_{r} \quad$ where
$\vec{r}^{\prime}=\vec{R}+\vec{r} \quad$ and
$\vec{r}^{\prime}$ refers to fixed(inertial system)
$\vec{r}$ refers to rotatinal(non-inertial system) rotates with $\vec{\omega}$ to $\vec{r}$ system
$\vec{R} \quad$ from the origin of $\vec{r}$ ' to the origin of $\vec{r}$
$\vec{v}_{r}=\left(\frac{d \vec{r}}{d t}\right)_{r}$

