UNIVERSITY OF SWAZILAND

## FACULTY OF SCIENCE

DEPARTMENT OF PHYSICS

SUPPLEMENTARY EXAMINATION 2016/2017

TITLE OF PAPER : CLASSICAL MECHANICS

COURSE NUMBER : P320

TIME ALLOWED : THREE HOURS

INSTRUCTIONS : ANSWER ANY FOUR OUT OF FIVE QUESTIONS.
EACH QUESTION CARRIES 25 MARKS.
MARKS FOR DIFFERENT SECTIONS ARE SHOWN IN THE RIGHT-HAND MARGIN.

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## Question one

(a) Given the following definite integral of $J(\alpha)=\int_{x_{1}}^{x_{2}} f\left(y(\alpha, x), y^{\prime}(\alpha, x), y^{\prime \prime}(\alpha, x) ; x\right) d x$, where the varied integration path is $y(\alpha, x)=y(x)+\alpha \eta(x), \eta\left(x_{1}\right)=\eta\left(x_{2}\right)=0$ and $\left.\frac{d \eta(x)}{d x}\right|_{x=x_{1}}=\left.\frac{d \eta(x)}{d x}\right|_{x=x_{2}}=0 \quad$ as shown in the following diagram :


Use the extremum condition for $J(\alpha)$, i.e., $\left.\frac{\partial J(\alpha)}{\partial \alpha}\right|_{\alpha=0}=0$, to deduce that
$f$ along the extremum path, i.e., $f\left(y(x), y^{\prime}(x), y^{\prime \prime}(x) ; x\right)$, satisfies the following equation:
$\frac{\partial f}{\partial y}-\frac{d}{d x}\left(\frac{\partial f}{\partial y^{\prime}}\right)+\frac{d^{2}}{d x^{2}}\left(\frac{\partial f}{\partial y^{\prime \prime}}\right)=0$.
( 12 marks)
(b) For a certain dynamical system the kinetic energy, $T$, and potential energy, $V$, are given by $T=\frac{1}{2}\left(\dot{q}_{1}^{2}+\dot{q}_{1} \dot{q}_{2}+\dot{q}_{2}^{2}\right)$ $V=\frac{3}{2} q_{2}^{2}$
where $q_{1}, q_{2}$ are the generalized coordinates.
(i) Write down Lagrange's equations of motion.
( 5 marks)
(ii) From the results in (b)(i), deduce a differential equation for $q_{2}$ only and show that its general solution can be written as $q_{2}(t)=k_{1} \cos (2 t)+k_{2} \sin (2 t)$ where $k_{1}$ and $k_{2}$ are arbitrary constants .
( 4 marks)
(iii) Substitute the general solution of $q_{2}$ into the results in (b)(i) and show that the the general solution of $q_{1}$ can be written as $q_{1}(t)=2 k_{1} \cos (2 t)+2 k_{2} \sin (2 t)+k_{3} t+k_{4}$ where $k_{3}$ and $k_{4}$ are arbitrary constants.

## Question two

A spherical pendulum of mass $m$ and length $b$ is shown in the figure below:

(a) (i) From $x=\frac{1}{b} \sin (\theta) \cos (\phi), y=b \sin (\theta) \sin (\phi) \& z=-b \cos (\theta)$ and $T=\frac{1}{2} m\left(\dot{x}^{2}+\dot{y}^{2}+\dot{z}^{2}\right) \& V=m g z$, deduce the following Lagrangian for the system in terms of $\theta \& \phi$ as
$L=\frac{1}{2} m b^{2}\left(\dot{\theta}^{2}+\dot{\phi}^{2} \sin ^{2}(\theta)\right)+m g b \cos (\theta) \quad \ldots \ldots$ (1)
(5 marks)
(ii) Write down the equations of motion for $\theta \& \phi$ and deduce that
$\begin{cases}\frac{d p_{\theta}}{d t}=m b^{2} \sin (\theta) \cos (\theta) \dot{\phi}^{2}-m g b \sin (\theta) & \cdots \cdots \\ \frac{d p_{\phi}}{d t}=0 & \ldots \ldots .\end{cases}$
where $p_{\theta}=m b^{2} \dot{\theta} \quad \& \quad p_{\phi}=m b^{2} \sin ^{2}(\theta) \dot{\phi}$
(iii) From eq.(3), one has $p_{\phi}=$ const. $\xrightarrow{\text { setas }} K$, deduce from eq.(2) the following equation for small $\theta$, i.e., $\left(\sin (\theta) \approx \theta\right.$ and $\cos (\theta) \approx 1-\frac{\theta^{2}}{2}$ or 1$)$, that $m^{2} b^{4} \theta^{3} \ddot{\theta}=K^{2}-m^{2} g b^{3} \theta^{4}$
(iv) If $K=0$ in eq.(4), write down the general solution of $\theta(t)$. ( 3 marks )
(b) (i) Find the Hamiltonian of the system in terms of $\theta, \phi, p_{\theta} \& p_{\phi}$.
(Hint : Since the Lagrangian in (a)(i) is not explicitly depending on $t$, thus the Hamiltonian $H$ equals to $T+V$ )
(ii) Write down the equations of motion for $H$ in (b)(i).

## Question three

(a) Given the Lagrangian for the two-body central force system as :
$L=T-V=\frac{1}{2} \mu\left(\dot{r}^{2}+r^{2} \dot{\theta}^{2}\right)+\frac{k}{r^{\prime \prime}}$
where $\mu$ is the reduced mass of the system, $k \& n$ are positive constants and $(r, \theta)$ are polar coordinates of the motion plane with its origin at the center of mass of the two-body system.
(i) Write down the Lagrange's equations of motion for the given system and show that

$$
\left\{\begin{array}{l}
\dot{\theta}=\frac{l}{\mu r^{2}} \\
\mu \ddot{r}-\frac{l^{2}}{\mu r^{3}}+\frac{n k}{r^{n+1}}=0
\end{array}\right.
$$

where the angular momentum $l$ is a constant.
( 5 marks)
(ii) In the case of circular orbits, i.e., $r=$ const, , use the results in (a)(i) and $T=\frac{1}{2} \mu\left(\dot{r}^{2}+r^{2} \dot{\theta}^{2}\right) \quad \& \quad V=\frac{k}{r^{n}}$ to find a relation between the kinetic and potential energies and show that $T=-\frac{n}{2} V$
( 5 marks)
(b) Starting from the law of conservation of angular momentum $l\left(=\mu r^{2} \dot{\theta}\right)$, derive Kepler's third law, i.e., the relation between the period $\tau$ of a closed orbit and the area $A$ of the closed orbit. Show that $\tau=\frac{2 \mu}{l} A$
(c) If an earth satellite of 400 kg mass is having a pure tangential speed $v_{\theta}(=r \dot{\theta})=10,000=10^{4} \mathrm{~m} / \mathrm{s}$ at its near-earth-point 500 km above the earth surface,
(i) calculate the values of $l$ and $E$ of this satellite,
(ii) calculate the values of the eccentricity, $\varepsilon$, and show that the orbit is an elliptical orbit. Also calculate its period.
(a) Two set of coordinate systems are having the same origins. The non-prime system (with position vector denoted as $\vec{r}$ and referred to as "rotating" system) is rotating with an angular velocity $\bar{\sigma}$ about the prime system (with position vector denoted as $\vec{r}^{\prime}$ and referred as "fixed" system and taken as an inertial system). Use the relation that
$\left(\frac{d \vec{F}}{d t}\right)_{\text {fxeed }}=\left(\frac{d \vec{F}}{d t}\right)_{\text {rotating }}+\bar{\omega} \times \vec{F} \quad$ for any vector field $\vec{F} \quad$ where
$\left(\frac{d \vec{F}}{d t}\right)_{\text {jixed }} \& \quad\left(\frac{d \vec{F}}{d t}\right)_{\text {routing }}$ are time derivatives of $\vec{F}$ with respect to the fixed and rotating systems respectively, deduce the following:
$\vec{a}_{e f f}=\vec{a}-\dot{\bar{\omega}} \times \vec{r}-\vec{\omega} \times(\vec{\omega} \times \vec{r})-2 \vec{\omega} \times \vec{v}_{r} \quad$ where
$\bar{v}_{f} \equiv\left(\frac{d \bar{r}^{\prime}}{d t}\right)_{\text {jixed }}, \vec{v}_{r} \equiv\left(\frac{d \bar{r}}{d t}\right)_{\text {roating }}, \vec{a} \equiv\left(\frac{d \bar{v}_{f}}{d t}\right)_{\text {fixed }}, \vec{a}_{e f f} \equiv\left(\frac{d \vec{v}_{r}}{d t}\right)_{\text {rotaing }}$
and $\quad \dot{\bar{\omega}} \equiv\left(\frac{d \vec{\omega}}{d t}\right)_{\text {roating }}$
(Hint : Starting with the equation $\vec{r}^{\prime}=\vec{r}$ (same origin) and taking the time derivative of the equation twice with respect to the fixed system.)
(b)


Referring to the diagram above and considering the body coordinate system ( $x, y, z$ ) has the same origin as the earth's fixed inertial system $\left(x^{\prime}, y^{\prime}, z^{\prime}\right)$, i.e., center of the earth . Hanging a motionless simple pendulum of length $L$ and mass $m$ near the earth surface at a northern latitude $\lambda$, the pendulum is supposed to pointing direct downward along $-\bar{e}_{=}$direction but instead it is pointing toward the ground with a small angular deviation of $\delta$ made with the true downward direction resulting from the centrifugal force $(-m \bar{\omega} \times(\bar{\omega} \times \bar{r}))$ as shown in the diagram below.

## Question four (continued)


(i) Show that $\delta \approx \frac{\omega^{2} r_{E} \cos (\lambda) \sin (\lambda)}{g-\omega^{2} r_{E} \cos ^{2}(\lambda)}$
( 12 marks )
(Hint: $\vec{F}_{e f f} \approx \vec{e}_{z}(-m g)-m \bar{\omega} \times(\vec{\omega} \times \vec{r}), \vec{r} \approx \vec{e}_{z}\left(r_{E}\right) \&$ $\left.\bar{\omega}=\vec{e}_{x}(-\omega \cos (\lambda))+\bar{e}_{z}(\omega \sin (\lambda))\right)$
(ii) Given $\lambda=30^{\circ}, L=10 \mathrm{~m}, r_{E}=6400 \mathrm{~km}, g=9.8 \mathrm{~m} / \mathrm{s}^{2}$ and $\omega=2 \pi \mathrm{rad} / \mathrm{day}=7.27 \times 10^{-5} \mathrm{rad} / \mathrm{s}$, determine the value of the deviation angle $\delta$ in terms of degree.
( 3 marks)

## Question five

Consider the motion of the bobs in the double pendulum system in the figure below.


Both pendulums are identical and having the length $b$ and bob of mass $m$. The motion of both bobs is restricted to lie in the plane of this paper, i.e., $x-y$ plane.
(i) For small $\theta_{1}$ and $\theta_{2}$, i.e., $\left(\sin (\theta) \approx \theta\right.$ and $\cos (\theta) \approx 1-\frac{\theta^{2}}{2}$ or 1$)$, show that the Lagrangian for the system can be expressed as:
$L=m b^{2} \dot{\theta}_{1}^{2}+\frac{1}{2} m b^{2} \dot{\theta}_{2}^{2}+m b^{2} \dot{\theta}_{1} \dot{\theta}_{2}-m g b\left(1+\theta_{1}^{2}+\frac{\theta_{2}^{2}}{2}\right)$
where the zero gravitational potential is set at the equilibrium position of the lower bob,
i.e., $\theta_{1}=0, \theta_{2}=0$ and $y=0$.
( 6 marks )
(ii) Write down the equations of motion and deduce that

$$
\left\{\begin{array}{c}
2 \ddot{\theta}_{1}+\ddot{\theta}_{2}=-2 \frac{g}{b} \theta_{1} \\
\ddot{\theta}_{1}+\ddot{\theta}_{2}=-\frac{g}{b} \theta_{2}
\end{array}\right.
$$

(iii) Deduce from the equations in (ii) the following :

$$
\left\{\begin{array}{l}
\ddot{\theta}_{1}=-2 \frac{g}{b} \theta_{1}+\frac{g}{b} \theta_{2} \\
\ddot{\theta}_{2}=2 \frac{g}{b} \theta_{1}-2 \frac{g}{b} \theta_{2}
\end{array}\right.
$$

## Question five (continued)

(iv) Set $\theta_{1}=\hat{X}_{1} e^{i \omega t}$ and $\theta_{2}=\hat{X}_{2} e^{i \omega t}$ (where $\hat{X}_{1}$ and $\hat{X}_{2}$ are constants) and deduce from the equations in (iii) the matrix equation $-\omega^{2} X=A X \quad$ where

$$
X=\binom{\hat{X}_{1}}{\hat{X}_{2}} \text { and } A=\left(\begin{array}{cc}
-\left(2 \frac{g}{b}\right) & \frac{g}{b} \\
2 \frac{g}{b} & -\left(2 \frac{g}{b}\right)
\end{array}\right)
$$

(v) Find the eigenfrequencies, $\omega$, of this coupled system.
(vi) Find the eigenvectors of this coupled system.

## Useful informations

$V=-\int \vec{F} \cdot d \vec{l}$ and reversely $\vec{F}=-\vec{\nabla} V$
$L=T-V=L\left(q_{1}, q_{2}, \cdots, q_{n}, \dot{q}_{1}, \dot{q}_{2}, \cdots, \dot{q}_{n}, t\right)$
$p_{\alpha}=\frac{\partial L}{\partial \dot{q}_{\alpha}} \quad$ and $\quad \dot{p}_{\alpha}=\frac{\partial L}{\partial q_{\alpha}}$
$H=\sum_{\alpha=1}^{n}\left(p_{\alpha} \dot{q}_{\alpha}\right)-L=H\left(q_{1}, q_{2}, \cdots, q_{n}, \dot{q}_{1}, \dot{q}_{2}, \cdots, \dot{q}_{n}, t\right)$
$\dot{q}_{\alpha}=\frac{\partial H}{\partial p_{\alpha}} \quad$ and $\quad \dot{p}_{\alpha}=-\frac{\partial H}{\partial q_{\alpha}}$
$[u, v] \equiv \sum_{\alpha=1}^{n}\left(\frac{\partial u}{\partial q_{\alpha}} \frac{\partial v}{\partial p_{\alpha}}-\frac{\partial u}{\partial p_{\alpha}} \frac{\partial v}{\partial q_{\alpha}}\right)$
$G=6.673 \times 10^{-11} \frac{\mathrm{Nm}^{2}}{\mathrm{~kg}^{2}}$
radius of earth $r_{E}=6.4 \times 10^{6} \mathrm{~m}$
mass of earth $m_{E}=6 \times 10^{24} \mathrm{~kg}$
earth attractive potential $\equiv-\frac{k}{r}$ where $k=G m m_{E}$
$\varepsilon=\sqrt{1+\frac{2 E l^{2}}{\mu k}} \quad\{(\varepsilon=0$, circle $),(0<\varepsilon<1$, ellipse $),(\varepsilon=1$, parabola $), \cdots\}$
$\mu=\frac{m_{1} m_{2}}{m_{1}+m_{2}} \approx m_{1}$ if $m_{2} \gg m_{1}$
For elliptical orbit,i.e., $0<\varepsilon<1$, then $\left\{\begin{array}{c}\text { semi-major } a=\frac{k}{2|E|} \\ \text { semi-minor } b=\frac{l}{\sqrt{2 \mu|E|}} \\ \text { period } \tau=\frac{2 \mu}{l}(\pi a b) \\ r_{\min }=a(1-\varepsilon) \& r_{\max }=a(1+\varepsilon)\end{array}\right.$
for plane polar $(r, \theta)$ system with unit vectors $\left(\bar{e}_{r}, \vec{e}_{\theta}\right)$, we have
$\left\{\begin{array}{l}\bar{v}=\vec{e}_{r} \dot{r}+\vec{e}_{\theta} r \dot{\theta} \\ \bar{a}=\bar{e}_{r}\left(\ddot{r}-r \dot{\theta}^{2}\right)+\vec{e}_{\theta}(2 \dot{r} \dot{\theta}+r \ddot{\theta})\end{array}\right.$
$\vec{\nabla} f=\vec{e}_{r} \frac{\partial f}{\partial r}+\vec{e}_{\theta} \frac{1}{r} \frac{\partial f}{\partial \theta}$
$I=\left(\begin{array}{ccc}\sum_{\alpha} m_{\alpha}\left(x_{\alpha, 2}^{2}+x_{\alpha, 3}^{2}\right) & -\sum_{\alpha} m_{\alpha} x_{\alpha, 1} x_{\alpha, 2} & -\sum_{\alpha} m_{\alpha} x_{\alpha, 1} x_{\alpha, 3} \\ -\sum_{\alpha} m_{\alpha} x_{\alpha, 2} x_{\alpha, 1} & \sum_{\alpha} m_{\alpha}\left(x_{\alpha, 1}^{2}+x_{\alpha, 3}^{2}\right) & -\sum_{\alpha} m_{\alpha} x_{\alpha, 2} x_{\alpha, 3} \\ -\sum_{\alpha} m_{\alpha} x_{\alpha, 3} x_{\alpha, 1} & -\sum_{\alpha} m_{\alpha} x_{\alpha, 3} x_{\alpha, 2} & \sum_{\alpha} m_{\alpha}\left(x_{\alpha, 1}^{2}+x_{\alpha, 2}^{2}\right)\end{array}\right)$
$\bar{F}_{e f f}=\bar{F}-m \ddot{\vec{R}}_{j}-m \dot{\bar{\omega}} \times \vec{r}-m \bar{\omega} \times(\bar{\omega} \times \vec{r})-2 m \bar{\omega} \times \vec{v}_{r} \quad$ where
$\vec{r}^{\prime}=\vec{R}+\vec{r} \quad$ and
$\vec{r}^{\prime}$ refers to fixed(inertial system)
$\vec{r}$ refers to rotatinal(non-inertial system) rotates with $\bar{\omega}$ to $\vec{r}^{\prime}$ system
$\bar{R}$ from the origin of $\bar{r}$ ' to the origin of $\vec{r}$
$\vec{v}_{r}=\left(\frac{d \vec{r}}{d t}\right)_{r}$

